

Quasi-normal modes and non-linearities

Béatrice Bonga - 22 March 2024

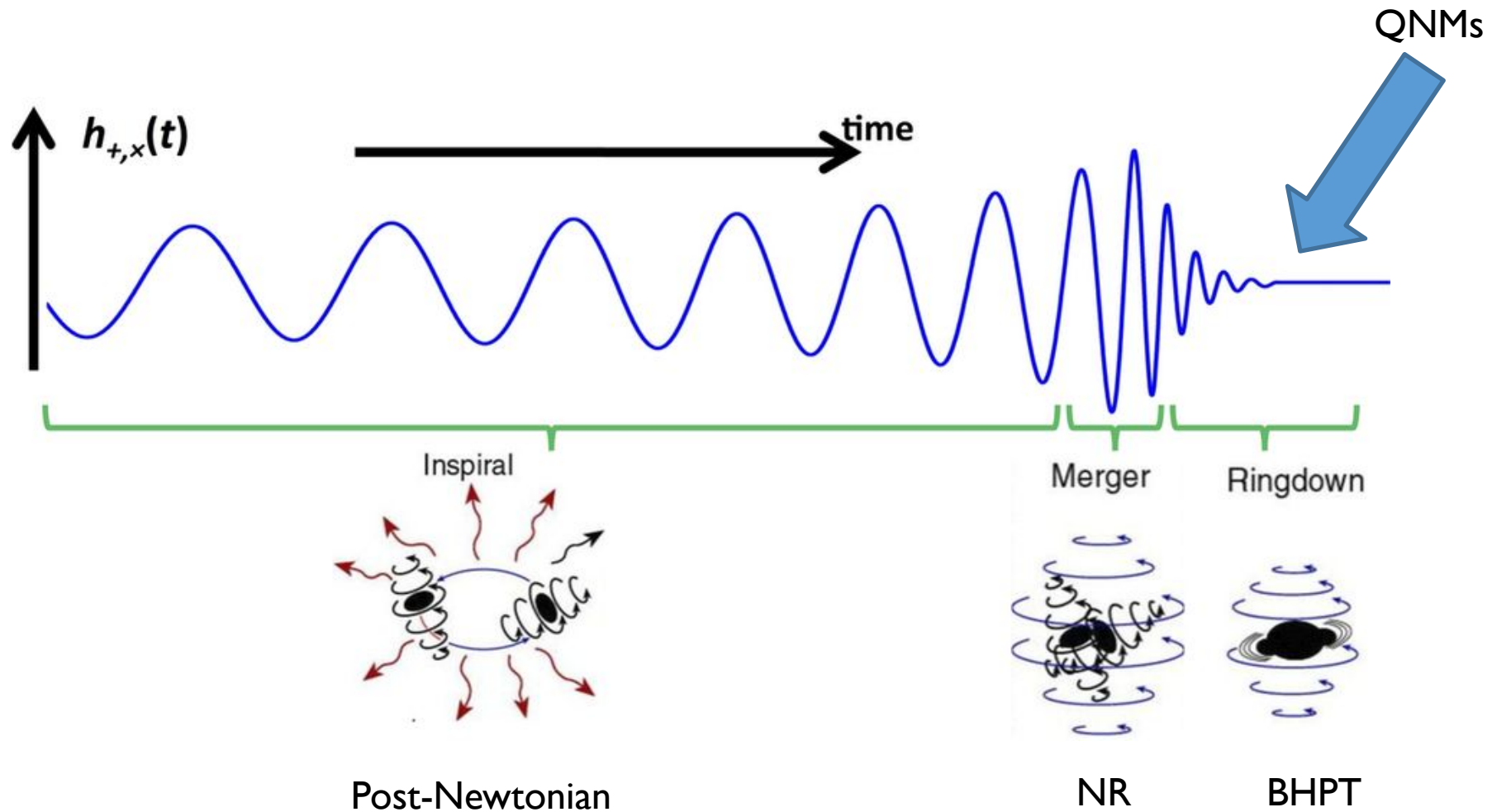
[Neev Khera, Ariadna Ribes Metidieri, BB, Xisco Jiménez Forteza, Badri
Krishnan, Eric Poisson, Daniel Pook-Kolb, Erik Schnetter, Huan Yang

PRL, arXiv:2306/11142]

Radboud University



Gravitational waves from black hole mergers



Quasi-normal modes



Mathematical description

$$h_+ + i h_\times = \sum_{l=2}^{\infty} \sum_{m=-l}^l h_{lm}(t, r) {}_2Y_{lm}(\theta, \phi)$$

spin-weighted
spherical harmonic

$$h_{lm}(t, r) = \frac{1}{r} \sum_{n=0}^N A_{lmn} e^{-i\omega_{lmn}(t-t_0) + \Phi_{lmn}}$$

Depend on the details
of the “hammer”

Frequencies and damping times

$$\omega_{lmn} = \omega_{lmn}^R + i \omega_{lmn}^I = 2\pi f_{lmn} + \frac{i}{\tau_{lmn}}$$



Depends on three integers:

$$\begin{aligned} l &= 2, 3, \dots \\ -l &< m < l \\ n &= 0, 1, 2, \dots \end{aligned}$$



Damping time

Frequencies can be calculated using black hole perturbation theory

Non-linearities?

Perturbation theory

$$\Psi^{(1)} \sim A_{\pm,lmn}^{(1)}(r) e^{-i\omega_{\pm,lmn}t + i\phi_{\pm,lmn}} {}_2Y_{lm}(\theta, \varphi)$$

$$\mathcal{O}\Psi^{(1)} = 0$$

$$\mathcal{O}\Psi^{(2)} = \mathcal{S}(h^{(1)}, h^{(1)})$$

$$\Psi^{(2)} \sim A_{\pm,lmn}^{(2)}(r) e_2^{-i\omega_{\pm,lmn}^{(2)}t + i\phi_{\pm,lmn}} Y_{lm}(\theta, \varphi)$$

$$\omega_{lmn \times l' m' n'} = \omega_{lmn} + \omega_{l' m' n'}$$

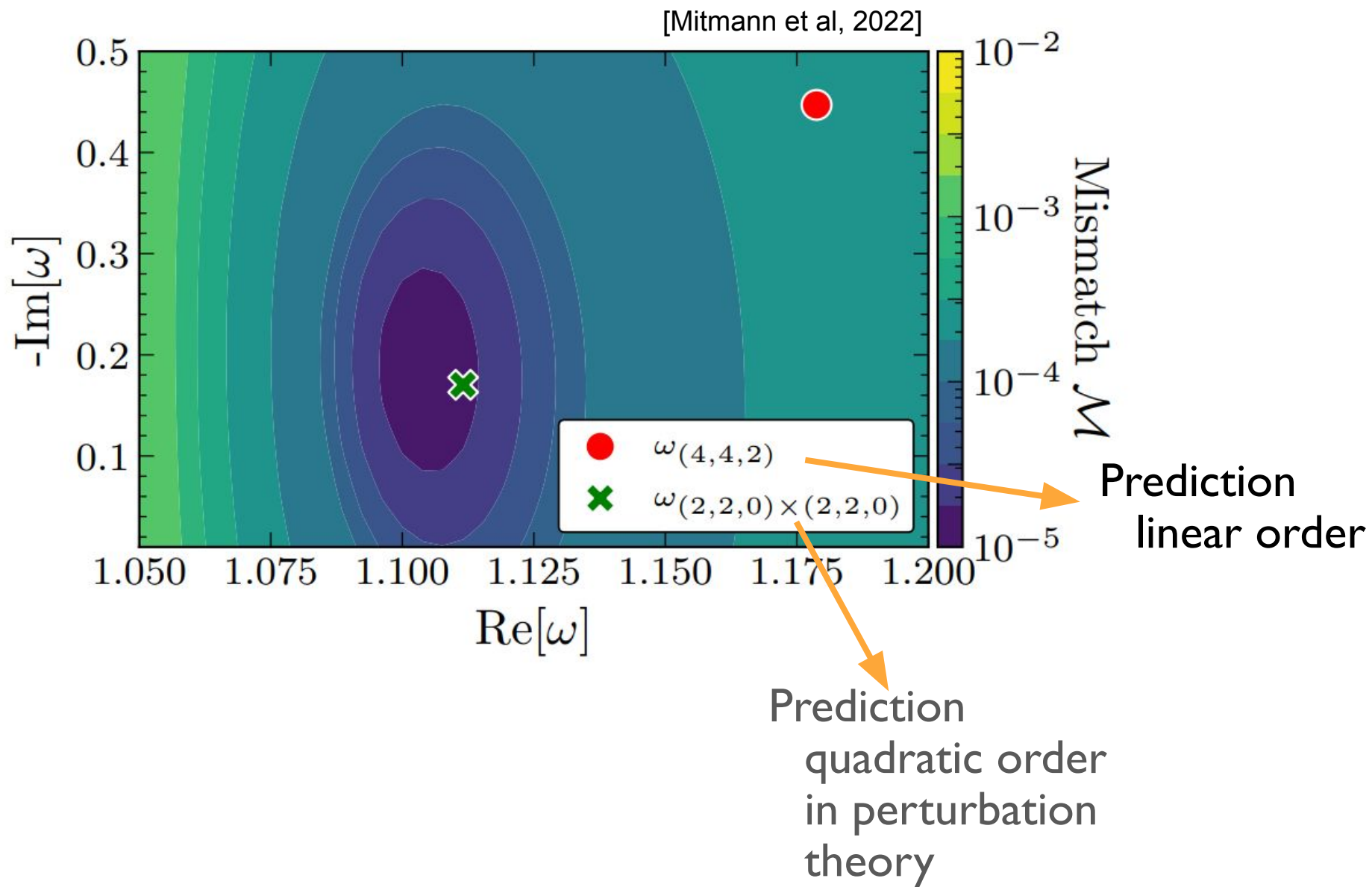
Amplitude relation

$$\mathcal{O}\Psi^{(2)} = \mathcal{S}(h^{(1)}, h^{(1)})$$

$$A_{lmn}^{(2)} Y_{lm} \sim \Sigma \overset{\text{background}}{\underbrace{f(r; M)}} \overset{\text{initial data}}{\underbrace{A_{lmn}^{(1)} A_{l'm'n'}^{(1)}}} \underbrace{Y_{lm} Y_{l'm'}}_{\sim G_{lm \times l'm'}} Y_{lm}$$

$$A_{lmn \times l'm'n'}^{(2)} = c_{lmn \times l'm'n'}(M, a) A_{lmn}^{(1)} A_{l'm'n'}^{(1)}$$

Non-linear model preferred @ infinity



So why do I think this is exciting?

Implications for observations:

$$h^{obs} = h^{linear} + h^{non-linear}$$

but frequencies are “finger-printed” with an order in perturbation theory!

Can we also model the black hole horizon with QNMs?

Horizon should be more non-linear, but not too crazy

→ easier to find quadratic QNMs

Horizon is strong field regime

→ hopeless to try to find any QNMs



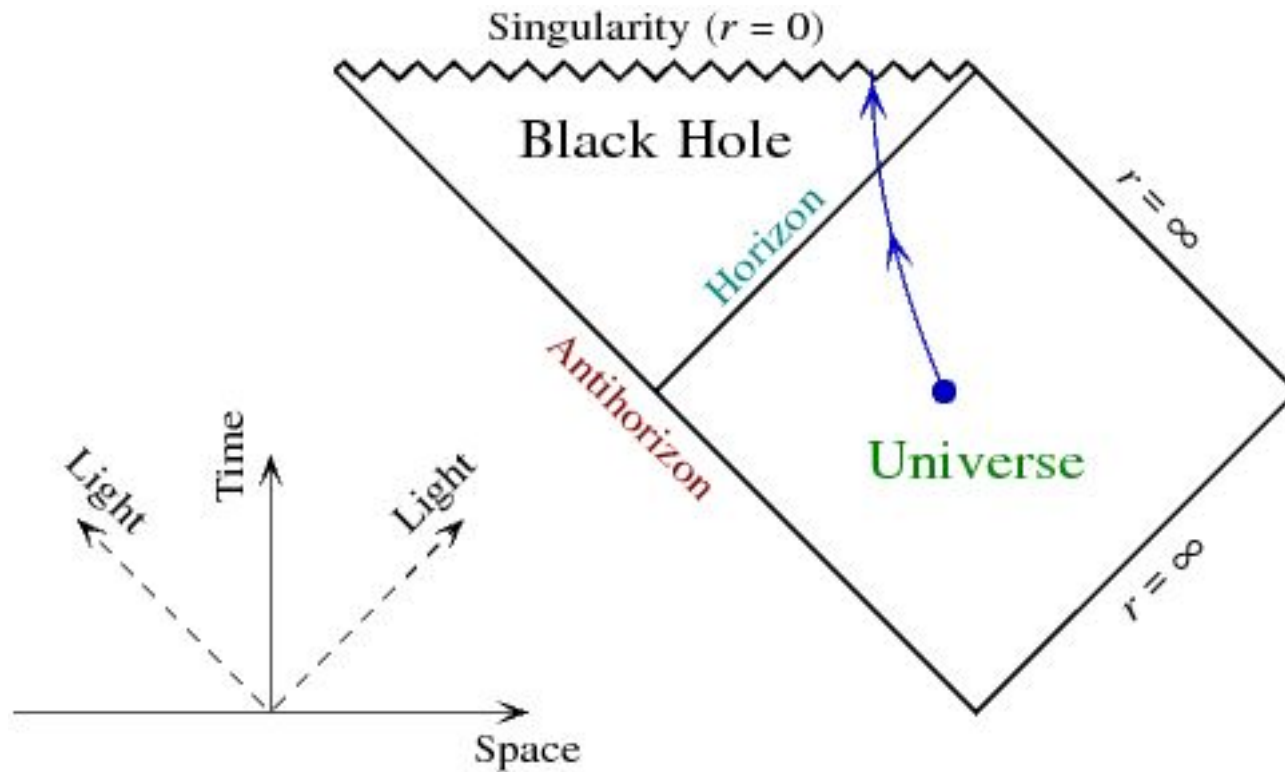
āngel



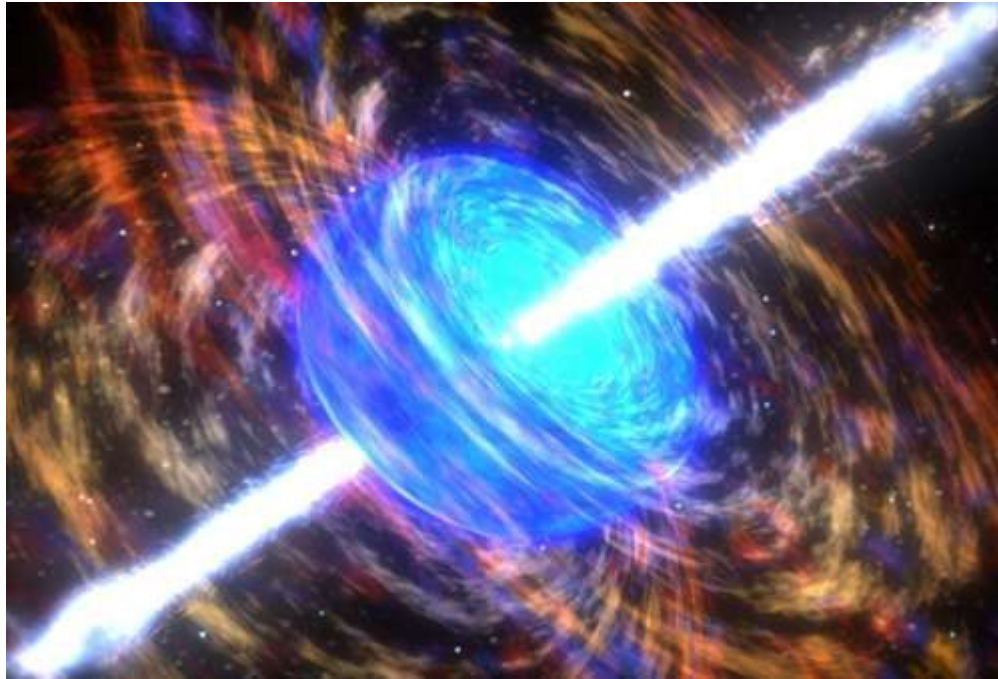
devil

Why care about the horizon...

...if observations are @ null infinity?



Electromagnetic observations and their sources



Gravitational waves...

are interesting because of their origin!

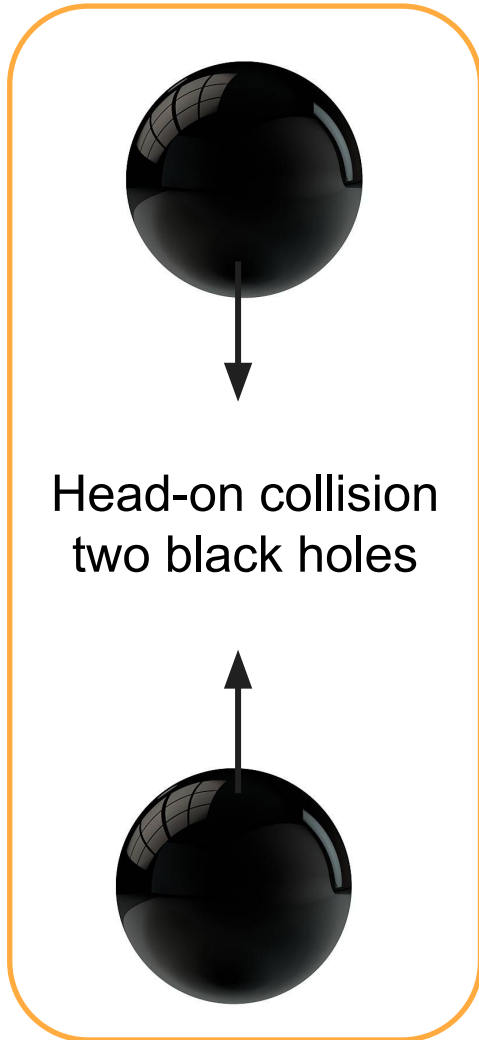
Corollary:

QNMs are interesting because they are emitted by black holes.

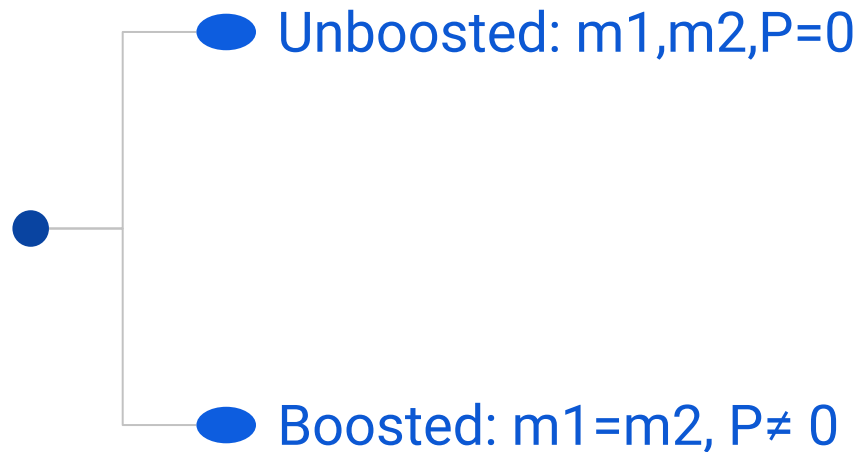
Disclaimer

*All results are based on fitting observations.
No theoretical derivations (yet)....*

Two sets of simulations using the Einstein Toolkit

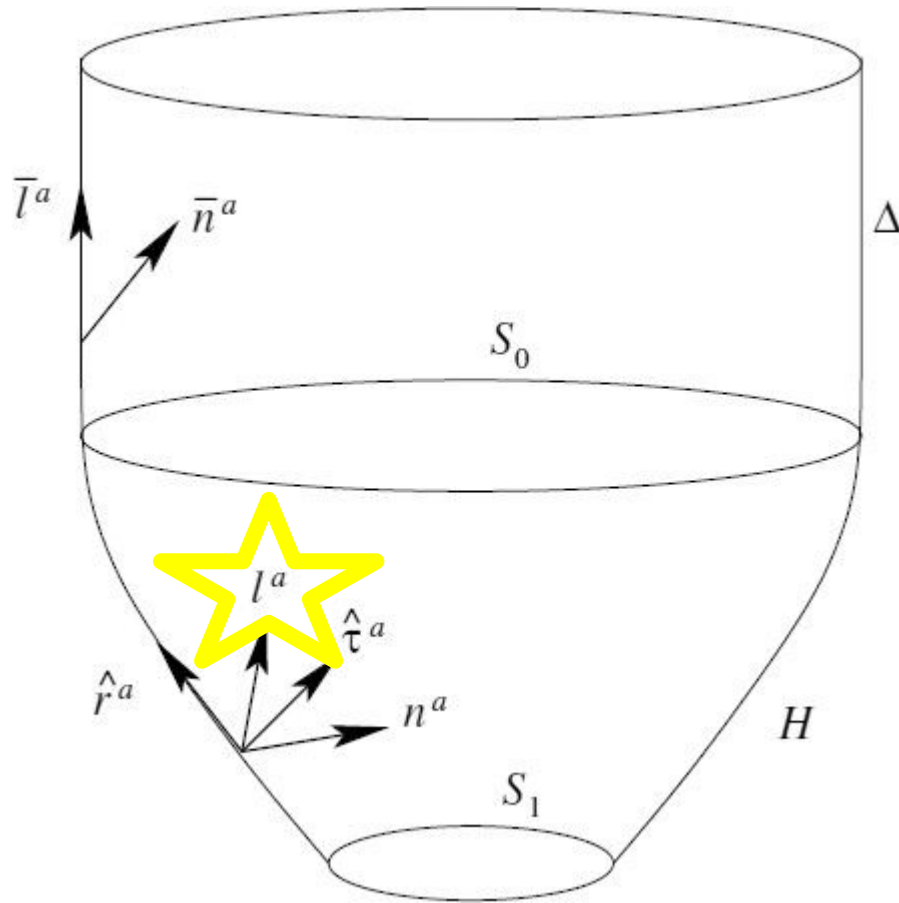


- (1) Resulting BH is non-rotating
- (2) Axisymmetric simulations \rightarrow no $m=0$ modes
- (3) High resolution near horizon (but poor near infinity)



 linear amplitudes 10x bigger

Shear at the horizon



Choice of time

Time

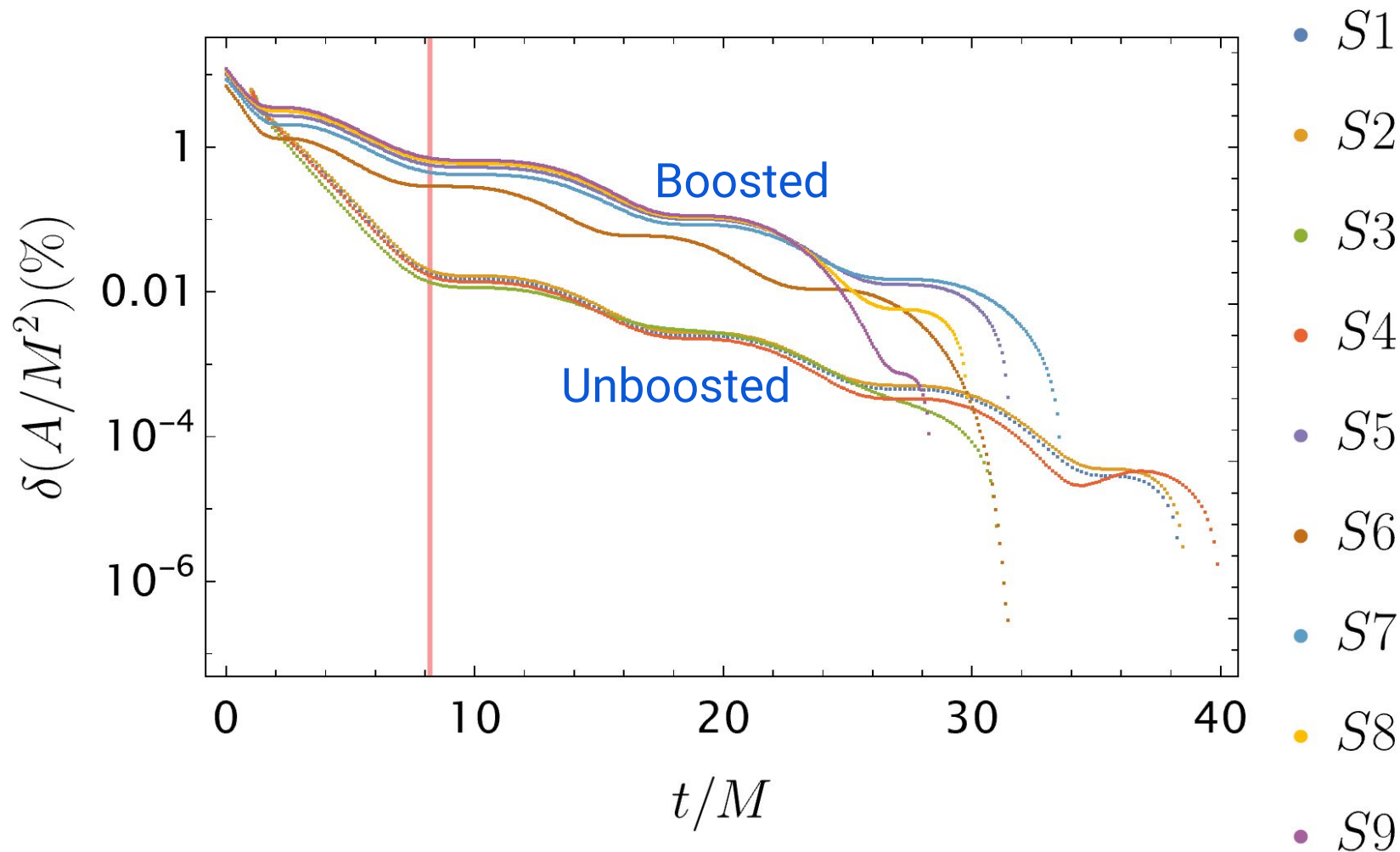


Definition of frequency

Disclaimer: We simply use the simulation time.

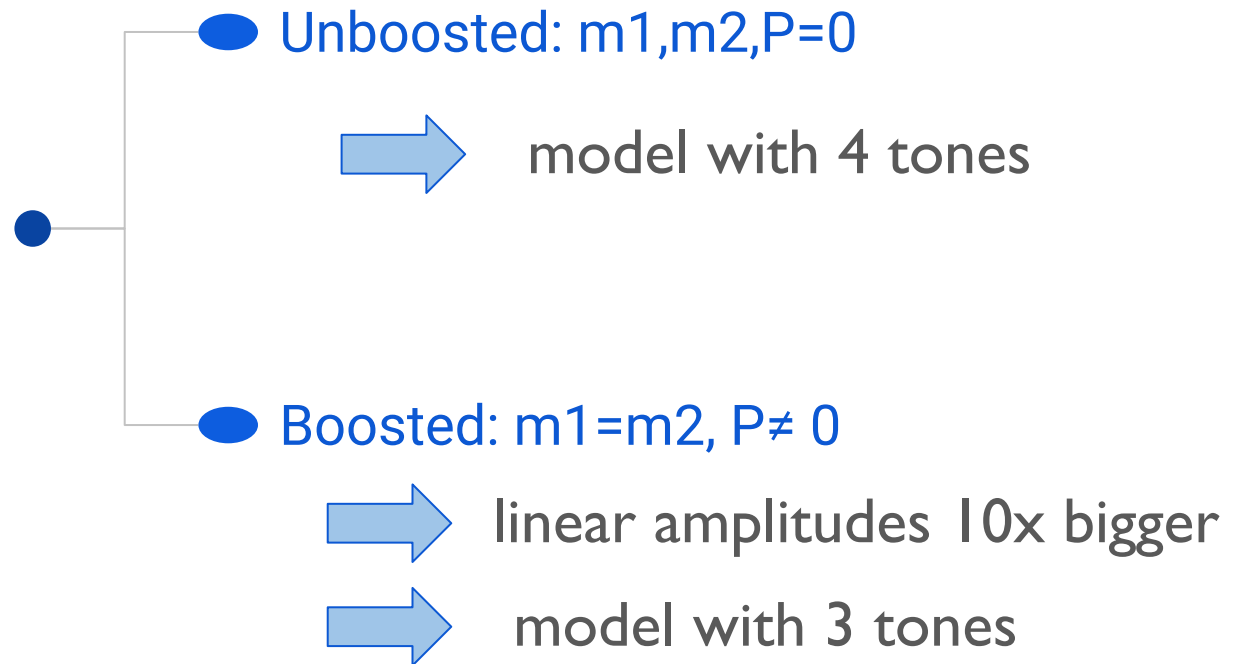
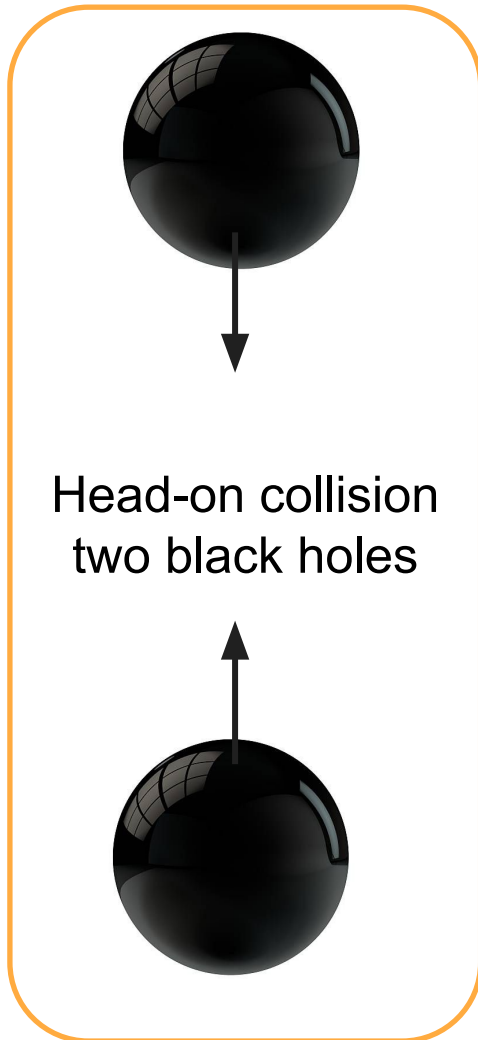
Same issue at infinity!

Ringdown: Mass changes $\leq 1\%$



We take $t_{\text{ringdown}} = 8.2 M$

Two sets of simulations using the Einstein Toolkit

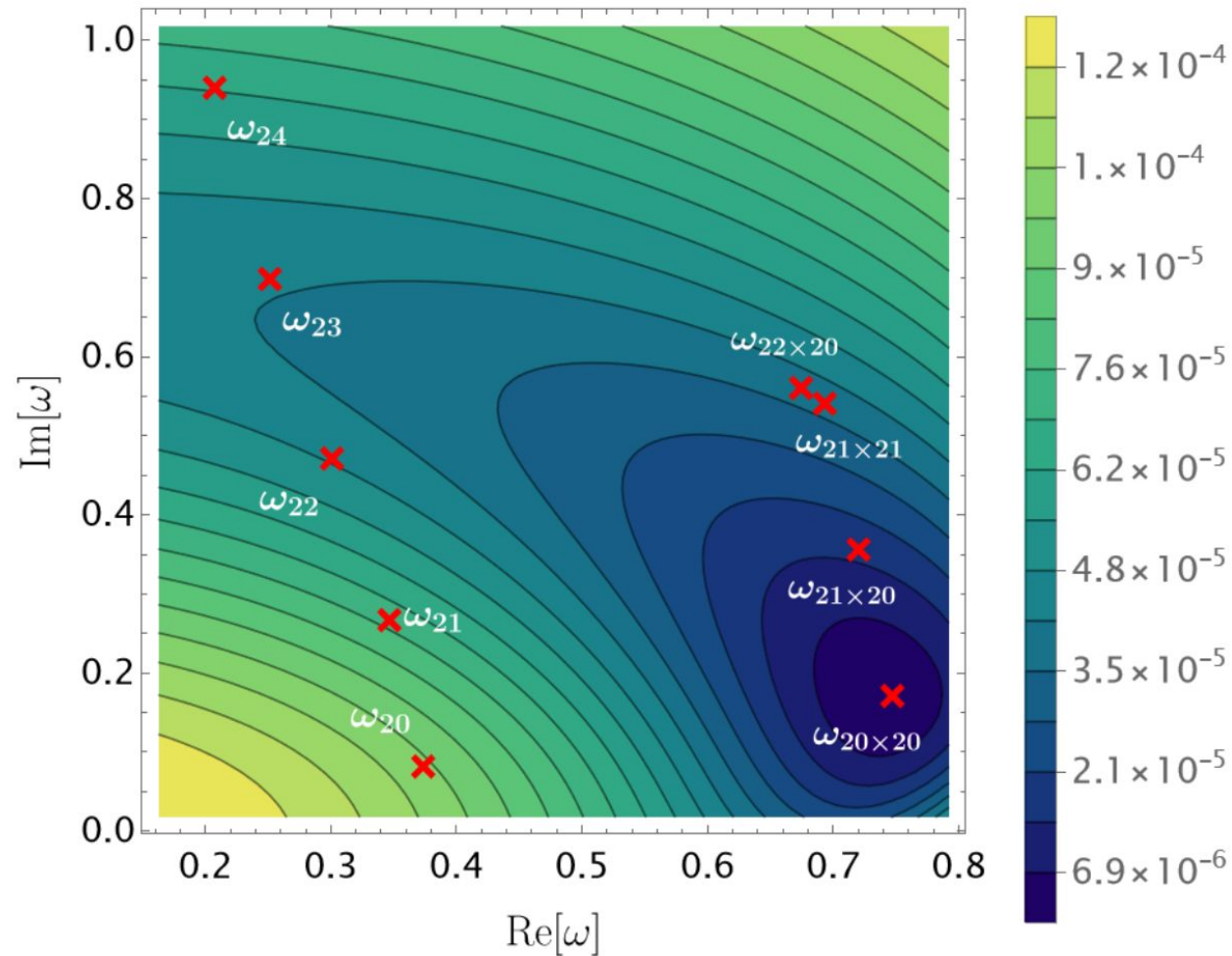


S7: boosted

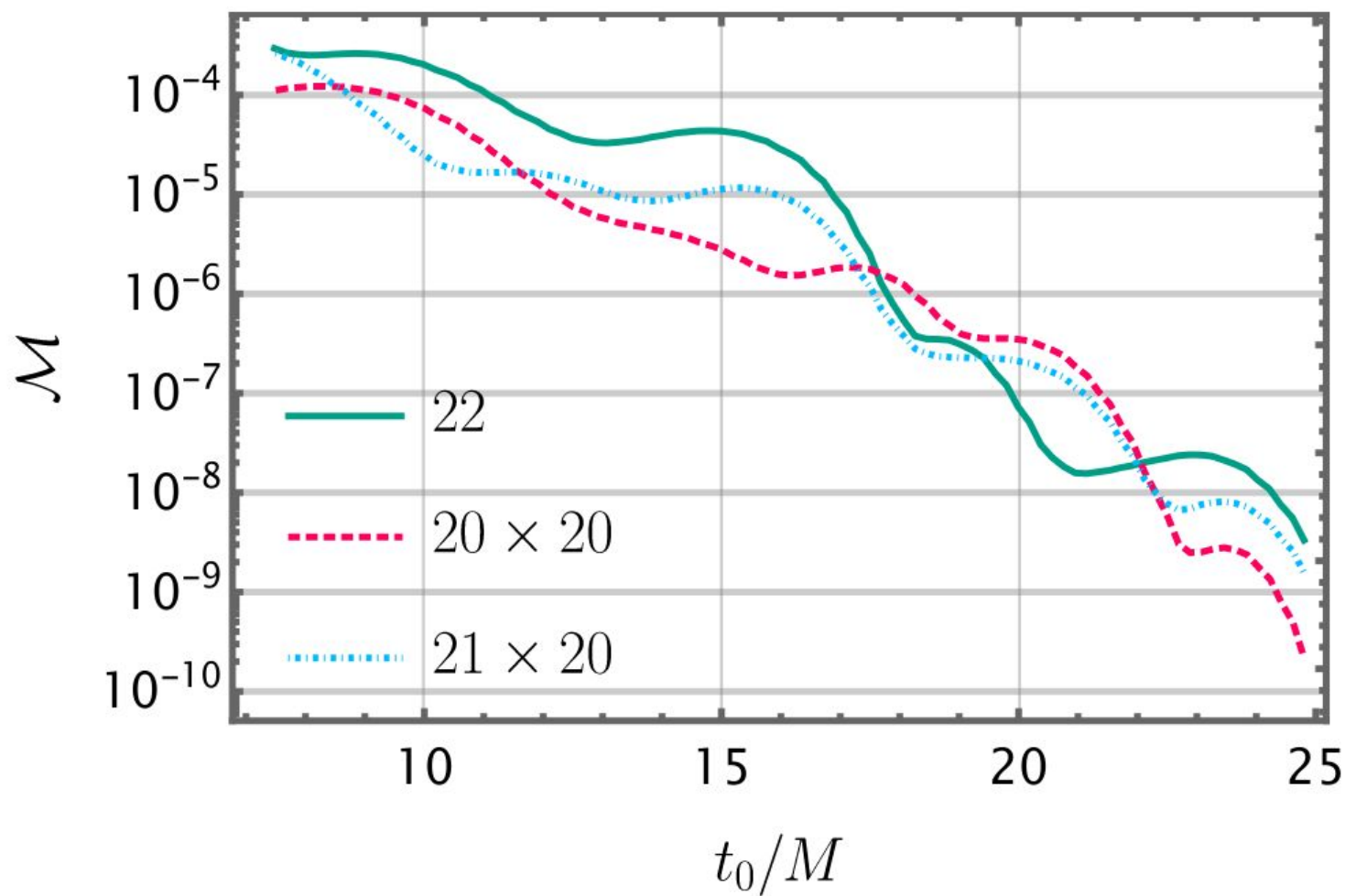
Equal mass \rightarrow $l=2,4,6,\dots$ are only non-zero.

Notation: $\omega_{lmn} \rightarrow \omega_{ln}$

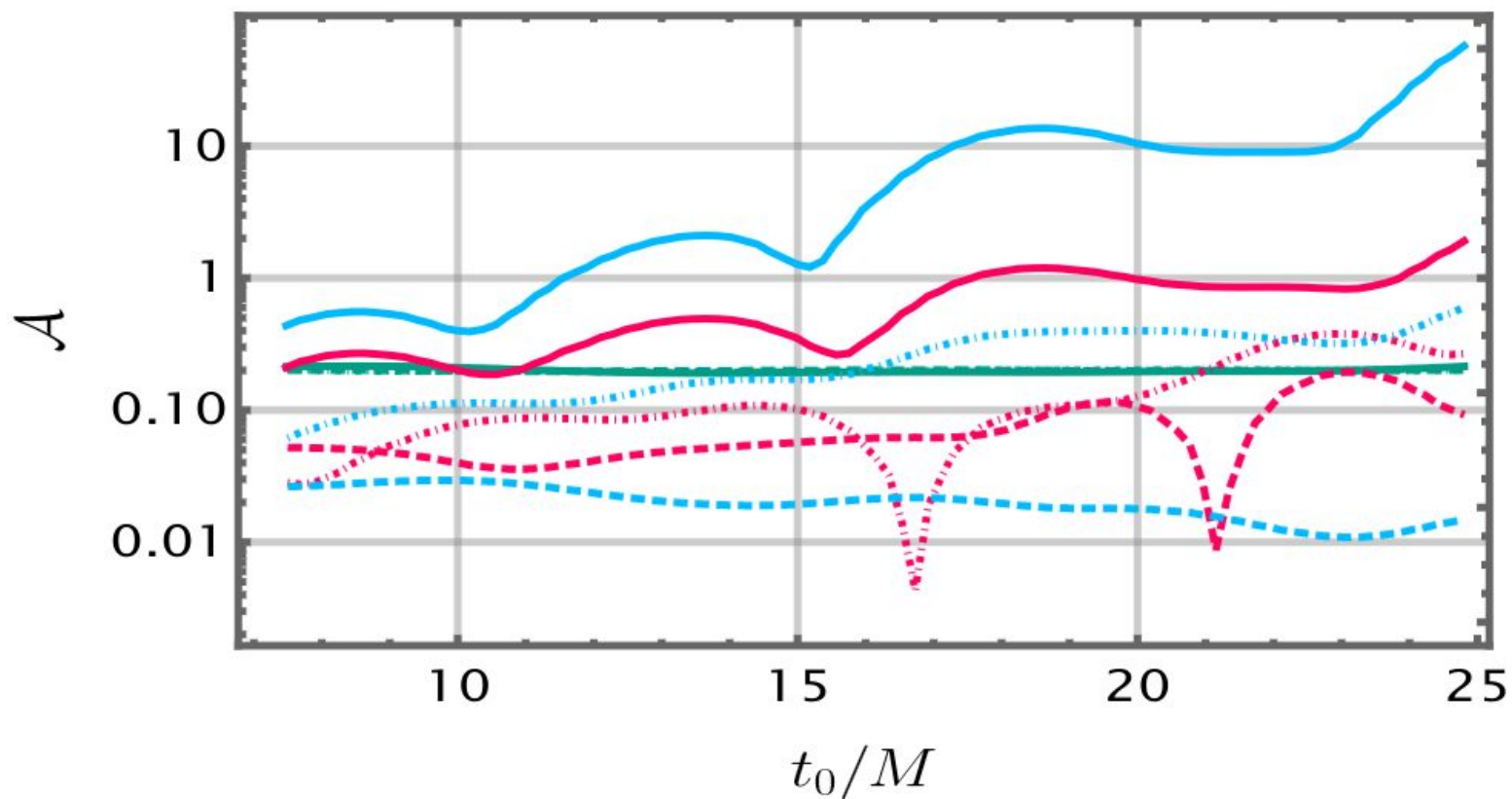
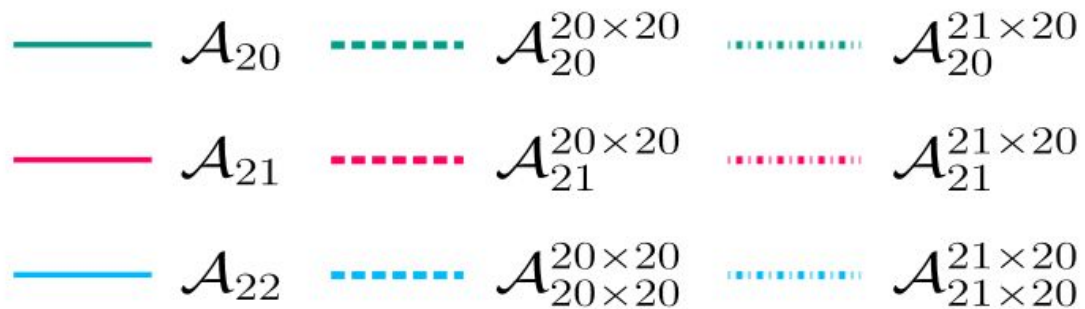
Mismatch S7 after fixing ω_{200} and ω_{201}



Mismatch S7 after fixing ω_{200} and ω_{201}

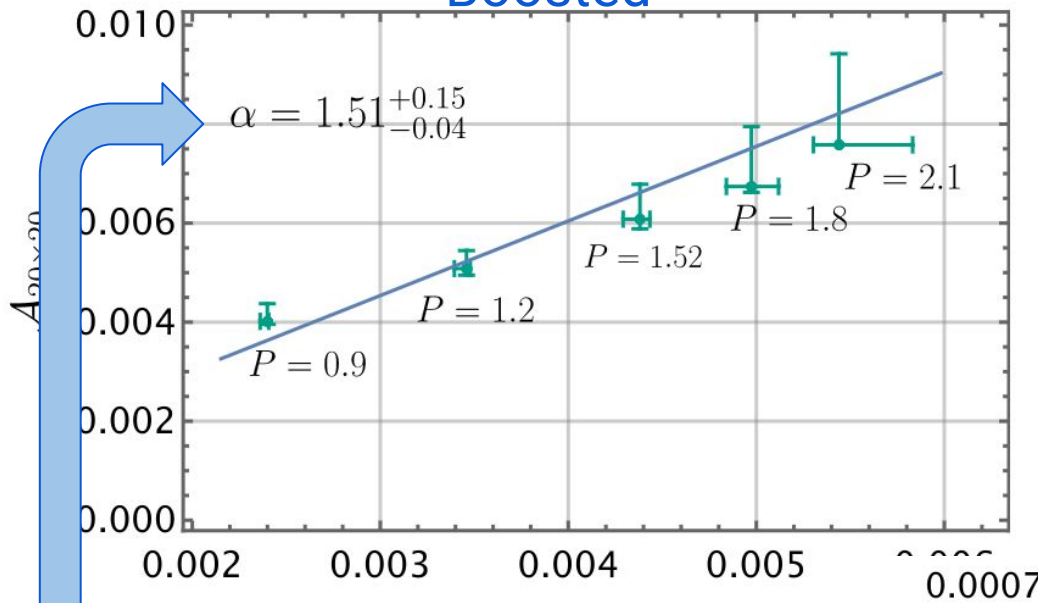


Stability amplitude

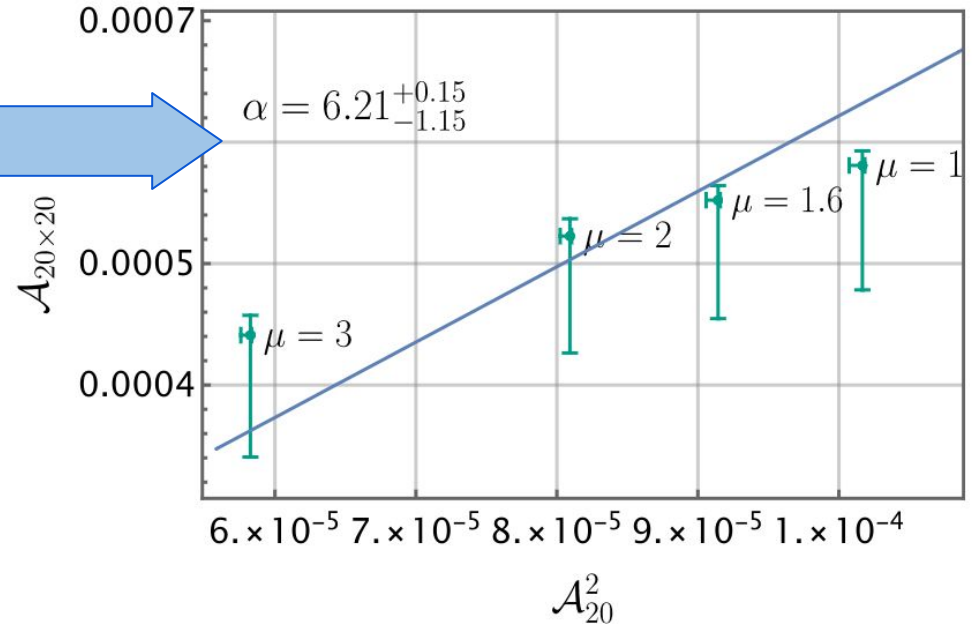


Amplitude relation

Boosted



Unboosted



Puzzle: Why are these slopes different?

Other l-modes

Mode	$\omega_{ln \times l'n'}$	Boosted (α)	Unboosted (α)
$l = 2$	$\omega_{20 \times 20}$	$1.51^{+0.15}_{-0.04}$	$6.21^{+0.15}_{-1.15}$
$l = 4$	$\omega_{20 \times 20}$	$0.73^{+0.06}_{-0.33}$	-
	$\omega_{20 \times 40}$	$2.6^{+0.26}_{-0.26}$	-
$l = 6^*$	$\omega_{20 \times 40}$	$1.78^{0.53}_{-0.74}$	-
	$\omega_{20 \times 60}$	$2.52^{+1.29}_{-0.59}$	-
	$\omega_{20 \times 40}$	$1.78^{0.44}_{-0.65}$	-
	$\omega_{40 \times 40}$	$2.82^{+1.5}_{-0.62}$	-

Connection horizon and infinity

- For $l=4$, same quadratic modes found at infinity
- For $l=6$, also $\omega_{200 \times 400}$ found at infinity

[Cheung et al, 2022 + private correspondence]

Conclusion

- ★ Quadratic QNMs fit the shear (and multipole) data at the horizon better than models with overtones
 - lower mismatch
 - more stable amplitudes wrt changes in starting time
 - closer to the optimal frequency
 - amplitude relation is satisfied

- ★ Some of the same (quadratic) modes found at horizon and infinity

- ★ Puzzling: why is the amplitude relation for boosted and unboosted simulations different?

Open questions

- (1) Why are the slopes for boosted/unboosted simulations different?
- (2) All results based on fitting observations, are there better ways to do this?
- (3) Is there a well-motivated choice of slicing/time?
- (4) Can we link observations at infinity more directly to horizon properties?